

Introduction to Quantum Information Processing

CS 467 / CS 667

Phys 467 / Phys 767

C&O 481 / C&O 681

Lecture 7 (2005)

Richard Cleve

DC 3524

cleve@cs.uwaterloo.ca

Course web site at:

<http://www.cs.uwaterloo.ca/~cleve>

Contents

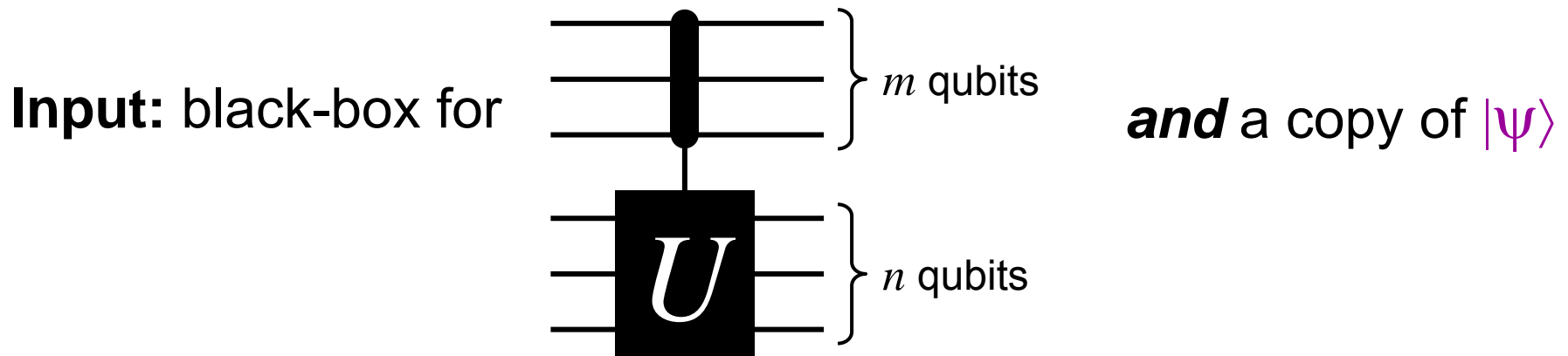
- Recap of phase estimation problem/algorithm
- How the algorithm works for general phases
- Recap of the order-finding problem/algorithm
- How to bypass the need for an eigenstate

- Recap of phase estimation problem/algorithm
- How the algorithm works for general phases
- Recap of the order-finding problem/algorithm
- How to bypass the need for an eigenstate

Phase estimation problem

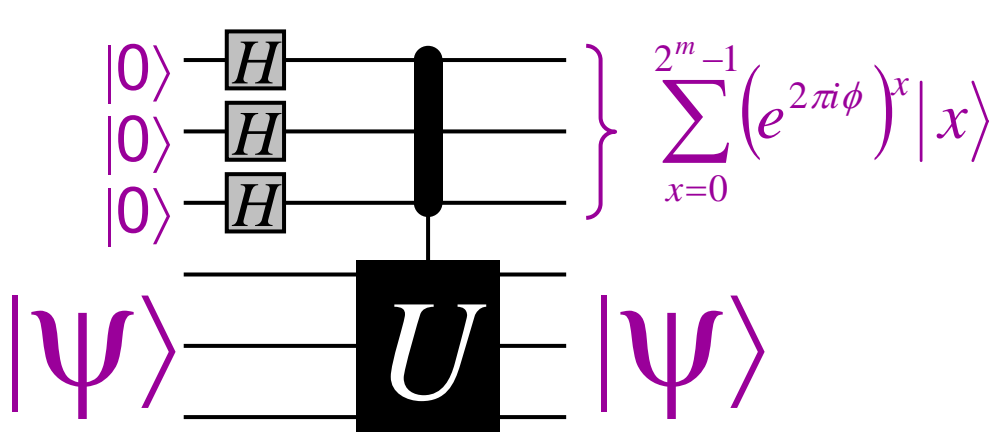
U is a unitary operation on n qubits

$|\psi\rangle$ is an eigenvector of U , with eigenvalue $e^{2\pi i\phi}$ ($0 \leq \phi < 1$)



Output: ϕ (m -bit approximation)

Algorithm for phase estimation



When $\phi = 0.a_1a_2\dots a_m$: $F_M |a_1a_2\dots a_m\rangle = \sum_{x=0}^{2^m-1} (e^{2\pi i \phi})^x |x\rangle$

$$F_M^{-1} = \frac{1}{\sqrt{M}} \begin{bmatrix} 1 & 1 & 1 & 1 & \dots & 1 \\ 1 & \omega^{-1} & \omega^{-2} & \omega^{-3} & \dots & \omega^{-(M-1)} \\ 1 & \omega^{-2} & \omega^{-4} & \omega^{-6} & \dots & \omega^{-2(M-1)} \\ 1 & \omega^{-3} & \omega^{-6} & \omega^{-9} & \dots & \omega^{-3(M-1)} \\ \vdots & \vdots & \vdots & \vdots & \ddots & \vdots \\ 1 & \omega^{-(M-1)} & \omega^{-2(M-1)} & \omega^{-3(M-1)} & \dots & \omega^{-(M-1)^2} \end{bmatrix}$$

Therefore, applying the **inverse** of F_M yields the digits of ϕ

- Recap of phase estimation problem/algorithm
- How the algorithm works for general phases
- Recap of the order-finding problem/algorithm
- How to bypass the need for an eigenstate

Arbitrary phases (I)

What if ϕ is not of the nice form $\phi = 0.a_1a_2\dots a_m$?

Example: $\phi = 1/3 = 0.\underline{0101010101010101}\dots$

Let's calculate what the previously-described procedure does:

Let $a/2^m = 0.a_1a_2\dots a_m$ be an m -bit approximation of ϕ ,
in the sense that $\phi = a/2^m + \delta$, where $|\delta| \leq 1/2^{m+1}$

$$\begin{aligned} (F_M)^{-1} \sum_{x=0}^{2^m-1} (e^{2\pi i \phi})^x |x\rangle &= \frac{1}{2^m} \sum_{y=0}^{2^m-1} \sum_{x=0}^{2^m-1} e^{-2\pi i xy/2^m} e^{2\pi i \phi x} |y\rangle \\ &= \frac{1}{2^m} \sum_{y=0}^{2^m-1} \sum_{x=0}^{2^m-1} e^{-2\pi i xy/2^m} e^{2\pi i \left(\frac{a}{2^m} + \delta\right) x} |y\rangle \\ &= \frac{1}{2^m} \sum_{y=0}^{2^m-1} \sum_{x=0}^{2^m-1} e^{2\pi i (a-y)x/2^m} e^{2\pi i \delta x} |y\rangle \end{aligned}$$

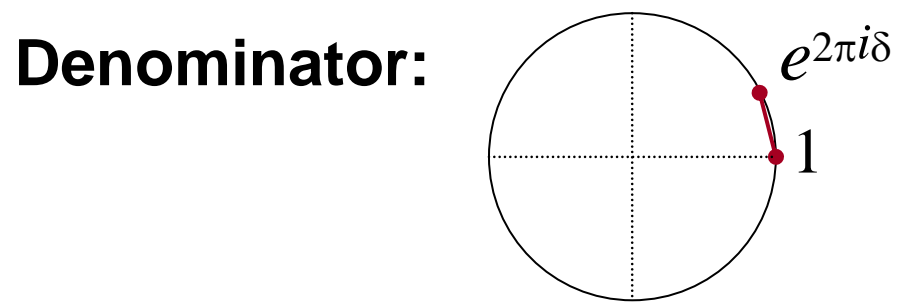
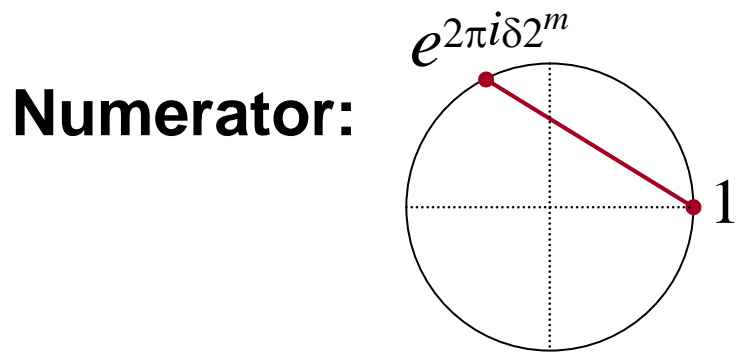
What is the
amplitude of
 $|a_1a_2\dots a_m\rangle$?

Arbitrary phases (II)

State is: $\frac{1}{2^m} \sum_{y=0}^{2^m-1} \sum_{x=0}^{2^m-1} e^{2\pi i(a-y)x/2^m} e^{2\pi i\delta x} |y\rangle$

geometric series!

The amplitude of $|y\rangle$, for $y = a$ is $\frac{1}{2^m} \sum_{x=0}^{2^m-1} e^{2\pi i\delta x} = \frac{1}{2^m} \frac{1 - (e^{2\pi i\delta})^{2^m}}{1 - e^{2\pi i\delta}}$



lower bounded by $2\pi i\delta 2^m (2/\pi) > 4\delta 2^m$

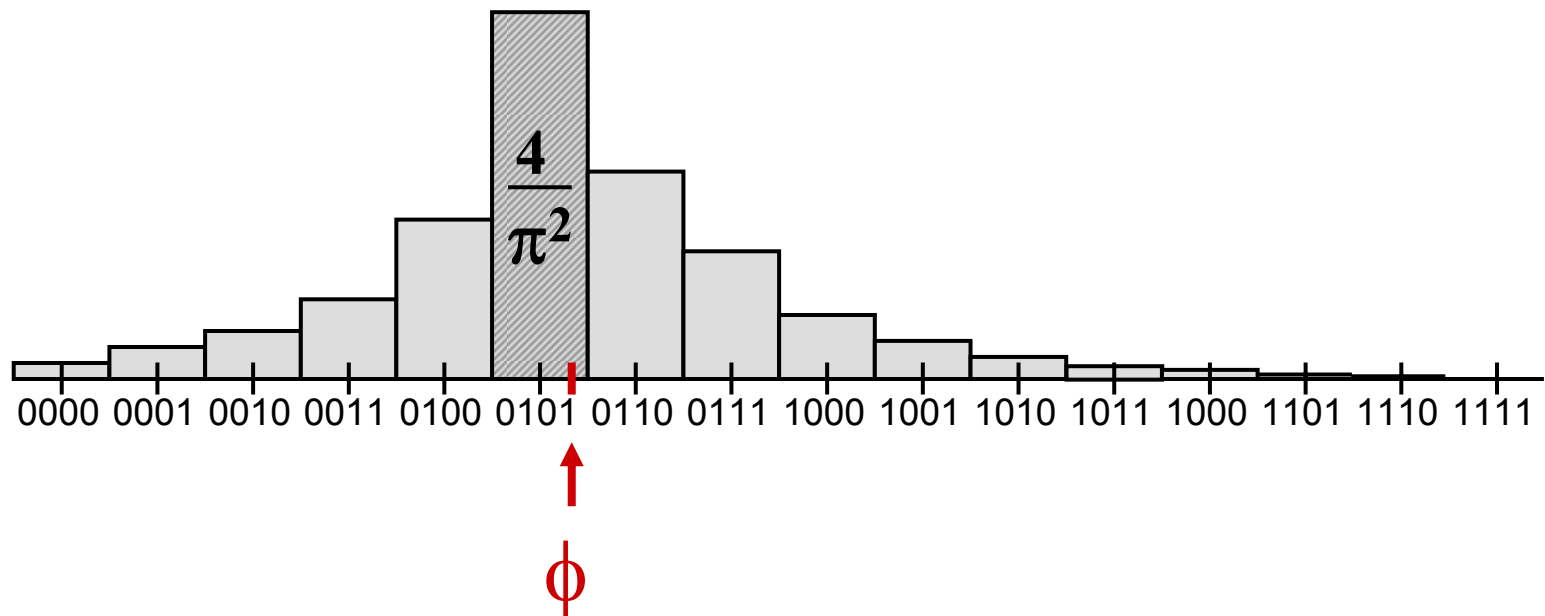
upper bounded by $2\pi\delta$

Therefore, the absolute value of the amplitude of $|y\rangle$ is at least the quotient of $(1/2^m)(\text{numerator}/\text{denominator})$, which is $2/\pi$

Arbitrary phases (III)

Therefore, the probability of measuring an m -bit approximation of ϕ is always at least $4/\pi^2 \approx 0.4$

For example, when $\phi = \frac{1}{3} = 0.\underline{01010101010101}\dots$, the outcome probabilities look roughly like this:



- Recap of phase estimation problem/algorithm
- How the algorithm works for general phases
- **Recap of the order-finding problem/algorithm**
- How to bypass the need for an eigenstate

Order-finding problem

Let M be an m -bit integer

Def: $\mathbf{Z}_M^* = \{x \in \{1, 2, \dots, M-1\} : \gcd(x, M) = 1\}$ (a group)

Def: $\text{ord}_M(a)$ is the minimum $r > 0$ such that $a^r = 1 \pmod{M}$

Order-finding problem: given a and M , find $\text{ord}_M(a)$

Example: $\mathbf{Z}_{21}^* = \{1, 2, 4, 5, 8, 10, 11, 13, 16, 17, 19, 20\}$

The powers of 10 are: 1, 10, 16, 13, 4, 19, 1, 10, 16, ...

Therefore, $\text{ord}_{21}(10) = 6$

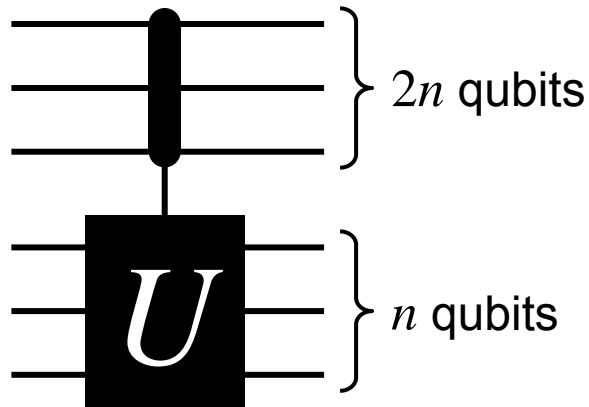
Order-finding algorithm (I)

Define: U (an operation on m qubits) as: $U|y\rangle = |ay \bmod M\rangle$

Define: $|\psi_1\rangle = \sum_{j=0}^{r-1} e^{-2\pi i(1/r)j} |a^j \bmod M\rangle$

Then $U|\psi_1\rangle = \sum_{j=0}^{r-1} e^{-2\pi i(1/r)j} |a^{j+1} \bmod M\rangle$
 $= \sum_{j=0}^{r-1} e^{2\pi i(1/r)} e^{-2\pi i(1/r)(j+1)} |a^{j+1} \bmod M\rangle$
 $= e^{2\pi i(1/r)} |\psi_1\rangle$

Order-finding algorithm (II)



corresponds to the mapping:

$$|x\rangle|y\rangle \rightarrow |x\rangle|a^x y \bmod M\rangle$$

Moreover, this mapping can be implemented with roughly $O(n^2)$ gates

The phase estimation algorithm yields a $2n$ -bit estimate of $1/r$

From this, a good estimate of r can be calculated by taking the reciprocal, and rounding off to the nearest integer

Exercise: why are $2n$ bits necessary and sufficient for this?

Problem: how do we construct state $|\psi_1\rangle$ to begin with?

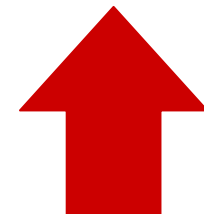
- Recap of phase estimation problem/algorithm
- How the algorithm works for general phases
- Recap of the order-finding problem/algorithm
- How to bypass the need for an eigenstate

Bypassing the need for $|\psi_1\rangle$ (I)

Let

$$\begin{aligned} |\psi_1\rangle &= \sum_{j=0}^{r-1} e^{-2\pi i(1/r)j} |a^j \bmod M\rangle \\ |\psi_2\rangle &= \sum_{j=0}^{r-1} e^{-2\pi i(2/r)j} |a^j \bmod M\rangle \\ &\vdots \\ |\psi_k\rangle &= \sum_{j=0}^{r-1} e^{-2\pi i(k/r)j} |a^j \bmod M\rangle \\ &\vdots \\ |\psi_r\rangle &= \sum_{j=0}^{r-1} e^{-2\pi i(r/r)j} |a^j \bmod M\rangle \end{aligned}$$

Can still uniquely determine k and r , provided they have no common factors (and $O(\log n)$ trials suffice for this)

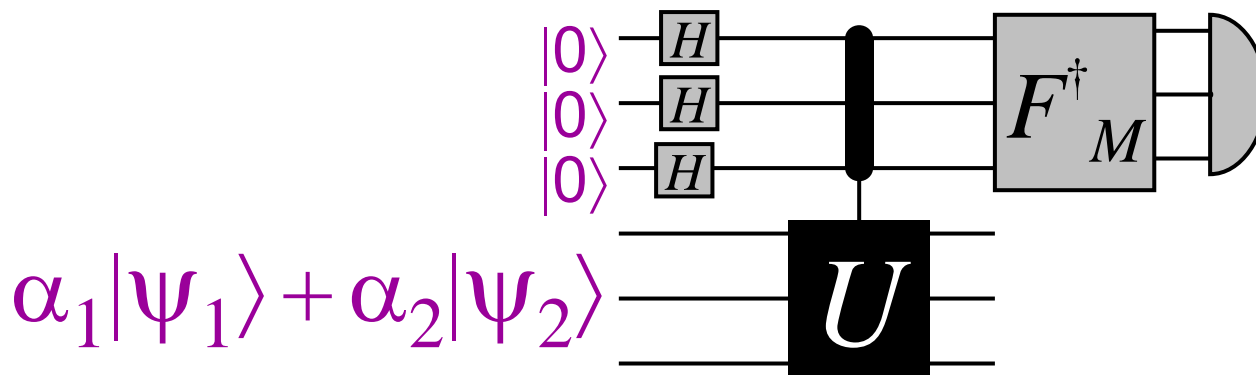


Any one of these could be used in the previous procedure, to yield an estimate of k/r , from which r can be extracted

What if k is chosen randomly and kept secret?

Bypassing the need for $|\psi_1\rangle$ (II)

Returning to the phase estimation problem, suppose that $|\psi_1\rangle$ and $|\psi_2\rangle$ have respective eigenvalues $e^{2\pi i\phi_1}$ and $e^{2\pi i\phi_2}$, and that $\alpha_1|\psi_1\rangle + \alpha_2|\psi_2\rangle$ is used in place of an eigenvalue:



What will the outcome be?

It will be an estimate of $\begin{cases} \phi_1 & \text{with probability } |\alpha_1|^2 \\ \phi_2 & \text{with probability } |\alpha_2|^2 \end{cases}$

Bypassing the need for $|\psi_1\rangle$ (III)

Using the state

yields results equivalent to choosing a $|\psi_k\rangle$ at random

Is it hard to construct the state $\frac{1}{\sqrt{r}} \sum_{k=1}^r |\psi_k\rangle$?

In fact, it's easy, since

$$\frac{1}{\sqrt{r}} \sum_{k=1}^r |\psi_k\rangle = \frac{1}{\sqrt{r}} \sum_{k=1}^r \sum_{j=0}^{r-1} e^{-2\pi i(k/r)j} |a^j \bmod M\rangle = |1\rangle$$

This is how the previous requirement for $|\psi_1\rangle$ is bypassed

THE END