Introduction to Quantum Information Processing

CS 467 / CS 667 Phys 667 / Phys 767

C&O 481 / C&O 681

Lecture 2 (2005)

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Superdense coding

How much classical information in n qubits?

 2^{n} —1 complex numbers apparently needed to describe an arbitrary n-qubit pure quantum state:

$$\alpha_{000}|000\rangle + \alpha_{001}|001\rangle + \alpha_{010}|010\rangle + ... + \alpha_{111}|111\rangle$$

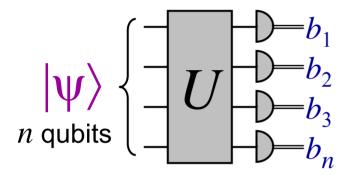
Does this mean that an exponential amount of classical information is somehow stored in n qubits?

Not in an operational sense ...

For example, Holevo's Theorem (from 1973) implies: one cannot convey more than n classical bits of information in n qubits

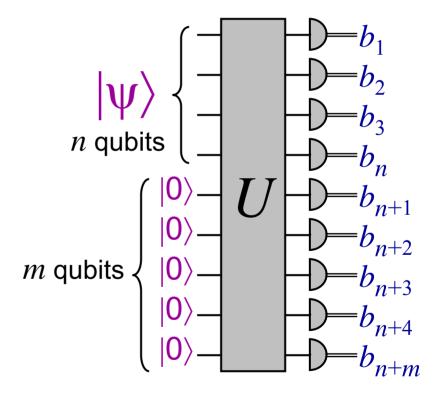
Holevo's Theorem

Easy case:



 $b_1b_2 \dots b_n$ certainly cannot convey more than n bits!

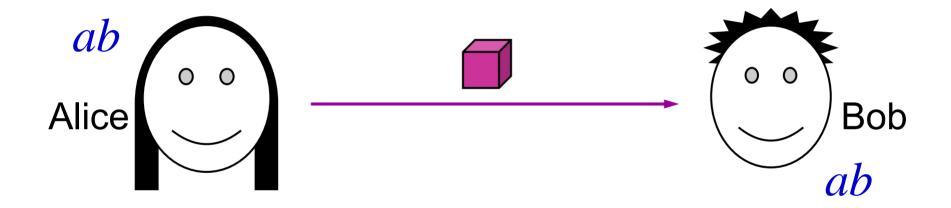
Hard case (the general case):



The difficult proof is beyond the scope of this course

Superdense coding (prelude)

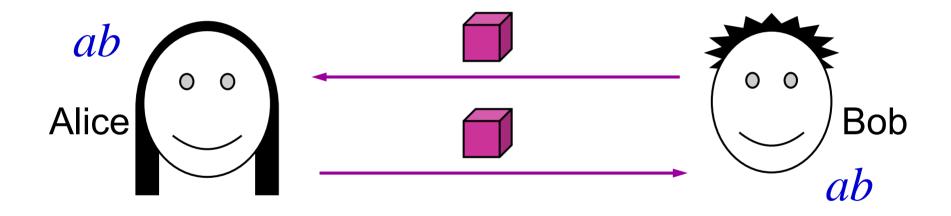
Suppose that Alice wants to convey **two** classical bits to Bob sending just **one** qubit



By Holevo's Theorem, this is *impossible*

Superdense coding

In *superdense coding*, Bob is allowed to send a qubit to Alice first



How can this help?

How superdense coding works

- 1. Bob creates the state $|00\rangle + |11\rangle$ and sends the *first* qubit to Alice
- 2. Alice: if a = 1 then apply X to qubit if b = 1 then apply Z to qubit send the qubit back to Bob

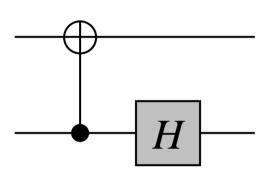
$$X = \begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix} \qquad Z = \begin{bmatrix} 1 & 0 \\ 0 & -1 \end{bmatrix}$$

ab	state	
00	00> + 11>	
01	00> - 11>	Bell basis
10	$ 01\rangle + 10\rangle$	
11	$ 01\rangle - 10\rangle$	

3. Bob measures the two qubits in the *Bell basis*

Measurement in the Bell basis

Specifically, Bob applies



input	output
$ 00\rangle + 11\rangle$	00⟩
$ 01\rangle + 10\rangle$	01⟩
00> - 11>	10⟩
01> - 10>	11⟩

to his two qubits ...

and then measures them, yielding *ab*

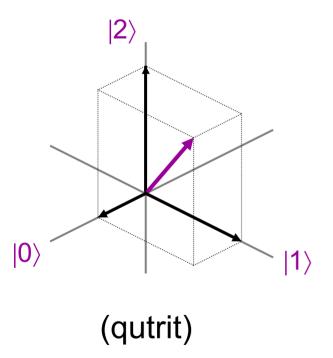
This concludes superdense coding

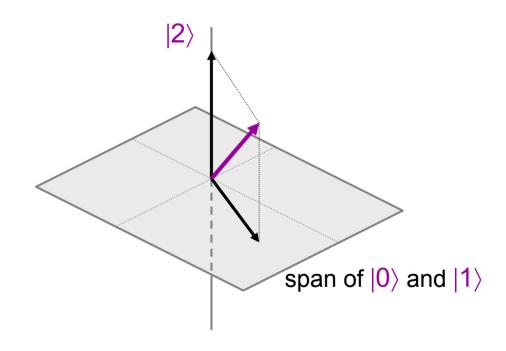
Teleportation

Incomplete measurements (I)

Measurements up until now are with respect to orthogonal one-dimensional subspaces:

The orthogonal subspaces can have other dimensions:





Incomplete measurements (II)

Such a measurement on $\alpha_0 | 0 \rangle + \alpha_1 | 1 \rangle + \alpha_2 | 2 \rangle$

$$(renormalized)$$
 results in
$$\begin{cases} \alpha_0|0\rangle+\alpha_1|1\rangle & \text{with prob } |\alpha_0|^2+|\alpha_1|^2\\ |2\rangle & \text{with prob } |\alpha_2|^2 \end{cases}$$

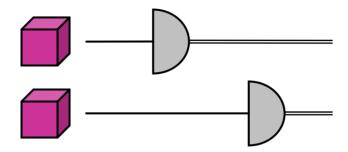
Measuring the first qubit of a two-qubit system

$$\alpha_{00}|00\rangle + \alpha_{01}|01\rangle + \alpha_{10}|10\rangle + \alpha_{11}|11\rangle \quad \left\{\begin{array}{c} \boxed{} \\ \boxed{} \end{array}\right.$$

Defined as the incomplete measurement with respect to the two dimensional subspaces:

- span of $|00\rangle$ & $|01\rangle$ (all states with first qubit 0), and
- span of $|10\rangle$ & $|11\rangle$ (all states with first qubit 1)

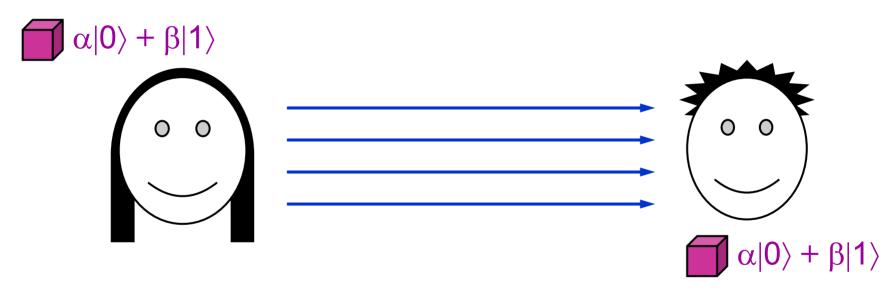
Result is the mixture
$$\begin{cases} \alpha_{00}|00\rangle+\alpha_{01}|01\rangle \text{ with prob } |\alpha_{00}|^2+|\alpha_{01}|^2\\ \alpha_{10}|10\rangle+\alpha_{11}|11\rangle \text{ with prob } |\alpha_{10}|^2+|\alpha_{11}|^2 \end{cases}$$



Easy exercise: show that measuring the first qubit and *then* measuring the second qubit gives the same result as measuring both qubits at once

Teleportation (prelude)

Suppose Alice wishes to convey a qubit to Bob by sending just classical bits

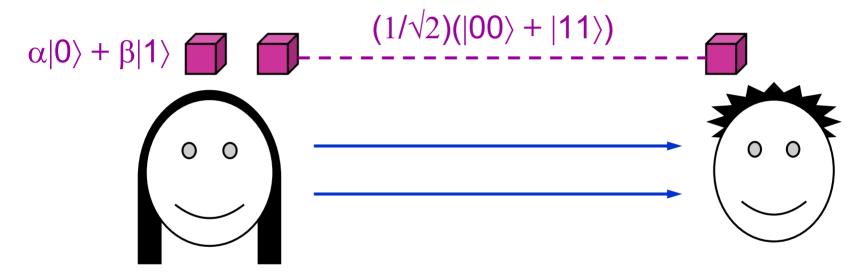


If Alice *knows* α and β , she can send approximations of them—but this still requires infinitely many bits for perfect precision

Moreover, if Alice does **not** know α or β , she can at best acquire **one bit** about them by a measurement

Teleportation scenario

In teleportation, Alice and Bob also start with a Bell state



and Alice can send two classical bits to Bob

Note that the initial state of the three qubit system is:

$$(1/\sqrt{2})(\alpha|0\rangle + \beta|1\rangle)(|00\rangle + |11\rangle)$$

= $(1/\sqrt{2})(\alpha|000\rangle + \alpha|011\rangle + \beta|100\rangle + \beta|111\rangle)$

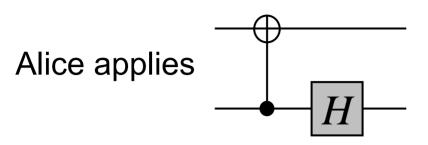
How teleportation works



Initial state:
$$(\alpha|0\rangle + \beta|1\rangle)(|00\rangle + |11\rangle)$$
 (omitting the $1/\sqrt{2}$ factor)
$$= \alpha|000\rangle + \alpha|011\rangle + \beta|100\rangle + \beta|111\rangle$$
$$= \frac{1}{2}(|00\rangle + |11\rangle)(\alpha|0\rangle + \beta|1\rangle)$$
$$+ \frac{1}{2}(|01\rangle + |10\rangle)(\alpha|1\rangle + \beta|0\rangle)$$
$$+ \frac{1}{2}(|00\rangle - |11\rangle)(\alpha|0\rangle - \beta|1\rangle)$$
$$+ \frac{1}{2}(|01\rangle - |10\rangle)(\alpha|1\rangle - \beta|0\rangle)$$

Protocol: Alice measures her two qubits *in the Bell basis* and sends the result to Bob (who then "corrects" his state) 16

What Alice does specifically



to her two qubits, yielding:

$$\begin{cases} \frac{1}{2}|00\rangle(\alpha|0\rangle+\beta|1\rangle) \\ +\frac{1}{2}|01\rangle(\alpha|1\rangle+\beta|0\rangle) \\ +\frac{1}{2}|10\rangle(\alpha|0\rangle-\beta|1\rangle) \\ +\frac{1}{2}|11\rangle(\alpha|1\rangle-\beta|0\rangle) \end{cases} \longrightarrow \begin{cases} (00,\alpha|0\rangle+\beta|1\rangle) & \text{with prob. } \frac{1}{4} \\ (01,\alpha|1\rangle+\beta|0\rangle) & \text{with prob. } \frac{1}{4} \\ (10,\alpha|0\rangle-\beta|1\rangle) & \text{with prob. } \frac{1}{4} \\ (11,\alpha|1\rangle-\beta|0\rangle) & \text{with prob. } \frac{1}{4} \end{cases}$$

Then Alice sends her two classical bits to Bob, who then adjusts his qubit to be $\alpha |0\rangle + \beta |1\rangle$ whatever case occurs

Bob's adjustment procedure

Bob receives two classical bits *a*, *b* from Alice, and:

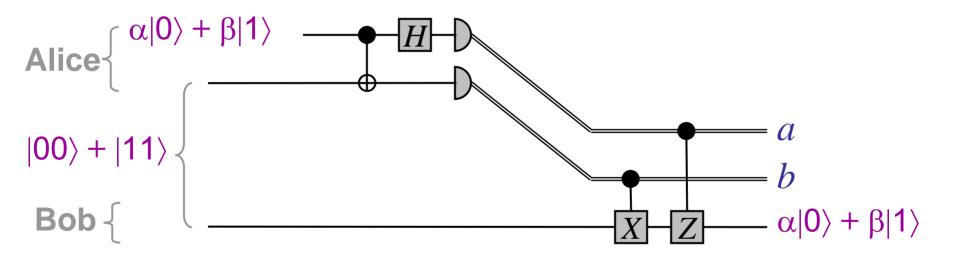
if
$$b = 1$$
 he applies X to qubit if $a = 1$ he applies Z to qubit

$$X = \begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix} \quad Z = \begin{bmatrix} 1 & 0 \\ 0 & -1 \end{bmatrix}$$

yielding:
$$\begin{cases} 00, & \alpha|0\rangle + \beta|1\rangle \\ 01, & X(\alpha|1\rangle + \beta|0\rangle) = \alpha|0\rangle + \beta|1\rangle \\ 10, & Z(\alpha|0\rangle - \beta|1\rangle) = \alpha|0\rangle + \beta|1\rangle \\ 11, & ZX(\alpha|1\rangle - \beta|0\rangle) = \alpha|0\rangle + \beta|1\rangle \end{cases}$$

Note that Bob acquires the correct state in each case

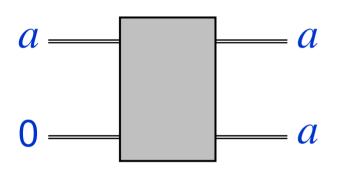
Summary of teleportation

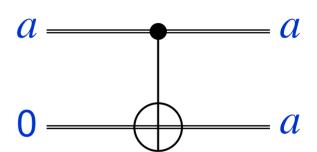


Quantum circuit exercise: try to work through the details of the analysis of this teleportation protocol

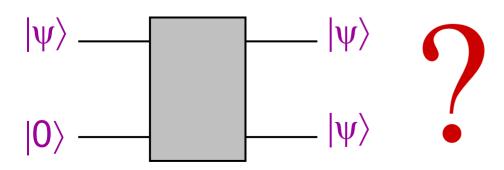
No-cloning theorem

Classical information can be copied

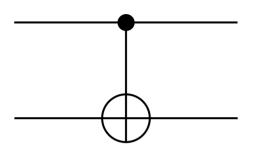




What about quantum information?



Candidate:



works fine for $|\psi\rangle = |0\rangle$ and $|\psi\rangle = |1\rangle$

... but it fails for $|\psi\rangle = (1/\sqrt{2})(|0\rangle + |1\rangle)$...

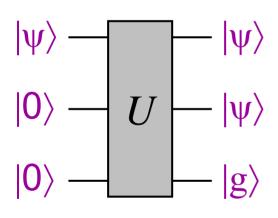
... where it yields output $(1/\sqrt{2})(|00\rangle + |11\rangle)$

instead of $|\psi\rangle|\psi\rangle = (1/4)(|00\rangle + |01\rangle + |10\rangle + |11\rangle)$

No-cloning theorem

Theorem: there is *no* valid quantum operation that maps an arbitrary state $|\psi\rangle$ to $|\psi\rangle|\psi\rangle$

Proof:



 $\begin{array}{c|c} |\psi\rangle & - & |\psi\rangle & \text{Let } |\psi\rangle \text{ and } |\psi'\rangle \text{ be two input s} \\ |0\rangle & - & |\psi\rangle & |\psi\rangle|\psi\rangle|g\rangle \text{ and } \\ |\psi\rangle|\psi\rangle|g\rangle \text{ and } \\ |\psi\rangle|\psi\rangle|g\rangle \text{ respectively} \\ |0\rangle & - & |g\rangle & \text{Since } U \text{ preserves inner product} \\ \end{array}$ Let $|\psi\rangle$ and $|\psi'\rangle$ be two input states,

Since U preserves inner products:

$$\langle \psi | \psi' \rangle = \langle \psi | \psi' \rangle \langle \psi | \psi' \rangle \langle g | g' \rangle \text{ so }$$
$$\langle \psi | \psi' \rangle (1 - \langle \psi | \psi' \rangle \langle g | g' \rangle) = 0 \text{ so }$$
$$|\langle \psi | \psi' \rangle| = 0 \text{ or } 1$$

