# Introduction to <br> Quantum Information Processing CS 467 I CS 667 Phys 667 I Phys 767 C\&O 481 / C\&O 681 

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## Superdense coding

## How much classical information in $\boldsymbol{n}$ qubits?

$2^{n}-1$ complex numbers apparently needed to describe an arbitrary $n$-qubit pure quantum state:
$\alpha_{000}|000\rangle+\alpha_{001}|001\rangle+\alpha_{010}|010\rangle+\ldots+\alpha_{111}|111\rangle$
Does this mean that an exponential amount of classical information is somehow stored in $n$ qubits?

## Not in an operational sense ...

For example, Holevo's Theorem (from 1973) implies: one cannot convey more than $n$ classical bits of information in $n$ qubits

## Holevo's Theorem

## Easy case:


$b_{1} b_{2} \ldots b_{n}$ certainly cannot convey more than $n$ bits!

Hard case (the general case):


The difficult proof is beyond the scope of this course

## Superdense coding (prelude)

Suppose that Alice wants to convey two classical bits to Bob sending just one qubit


By Holevo's Theorem, this is impossible

## Superdense coding

In superdense coding, Bob is allowed to send a qubit to Alice first


How can this help?

## How superdense coding works

1. Bob creates the state $|00\rangle+|11\rangle$ and sends the first qubit to Alice
2. Alice: if $a=1$ then apply $X$ to qubit if $b=1$ then apply $Z$ to qubit

$$
X=\left[\begin{array}{ll}
0 & 1 \\
1 & 0
\end{array}\right] \quad Z=\left[\begin{array}{rr}
1 & 0 \\
0 & -1
\end{array}\right]
$$ send the qubit back to Bob

$\left.\begin{array}{|l|l|}\hline a b & \text { state } \\ \hline 00 & |00\rangle+|11\rangle \\ 01 & |00\rangle-|11\rangle \\ 10 & |01\rangle+|10\rangle \\ 11 & |01\rangle-|10\rangle \\ \hline\end{array}\right\}$ Bell basis
3. Bob measures the two qubits in the Bell basis

## Measurement in the Bell basis

Specifically, Bob applies


| input | output |
| :--- | :---: |
| $\|00\rangle+\|11\rangle$ | $\|00\rangle$ |
| $\|01\rangle+\|10\rangle$ | $\|01\rangle$ |
| $\|00\rangle-\|11\rangle$ | $\|10\rangle$ |
| $\|01\rangle-\|10\rangle$ | $\|11\rangle$ |

to his two qubits ... and then measures them, yielding $a b$

This concludes superdense coding

## Teleportation

## Incomplete measurements (I)

Measurements up until now are with respect to orthogonal one-dimensional subspaces:


The orthogonal subspaces can have other dimensions:


## Incomplete measurements (II)

Such a measurement on $\alpha_{0}|0\rangle+\alpha_{1}|1\rangle+\alpha_{2}|2\rangle$
(renormalized)
results in $\left\{\begin{array}{cl}\alpha_{0}|0\rangle+\alpha_{1}|1\rangle & \text { with prob }\left|\alpha_{0}\right|^{2}+\left|\alpha_{1}\right|^{2} \\ |2\rangle & \text { with prob }\left|\alpha_{2}\right|^{2}\end{array}\right.$

## Measuring the first qubit of a two-qubit system

$$
\alpha_{00}|00\rangle+\alpha_{01}|01\rangle+\alpha_{10}|10\rangle+\alpha_{11}|11\rangle
$$



Defined as the incomplete measurement with respect to the two dimensional subspaces:

- span of $|00\rangle \&|01\rangle$ (all states with first qubit 0), and
- span of $|10\rangle \&|11\rangle$ (all states with first qubit 1)

Result is the mixture $\left\{\begin{array}{l}\alpha_{00}|00\rangle+\alpha_{01}|01\rangle \text { with prob }\left|\alpha_{00}\right|^{2}+\left|\alpha_{01}\right|^{2} \\ \alpha_{10}|10\rangle+\alpha_{11}|11\rangle \text { with prob }\left|\alpha_{10}\right|^{2}+\left|\alpha_{11}\right|^{2}\end{array}\right.$


Easy exercise: show that measuring the first qubit and then measuring the second qubit gives the same result as measuring both qubits at once

## Teleportation (prelude)

Suppose Alice wishes to convey a qubit to Bob by sending just classical bits
$\square \alpha|0\rangle+\beta|1\rangle$


If Alice knows $\alpha$ and $\beta$, she can send approximations of them -but this still requires infinitely many bits for perfect precision

Moreover, if Alice does not know $\alpha$ or $\beta$, she can at best acquire one bit about them by a measurement

## Teleportation scenario

In teleportation, Alice and Bob also start with a Bell state

and Alice can send two classical bits to Bob
Note that the initial state of the three qubit system is:

$$
\begin{aligned}
& (1 / \sqrt{ } 2)(\alpha|0\rangle+\beta|1\rangle)(|00\rangle+|11\rangle) \\
& =(1 / \sqrt{ } 2)(\alpha|000\rangle+\alpha|011\rangle+\beta|100\rangle+\beta|111\rangle)
\end{aligned}
$$

## How teleportation works <br> 

Initial state: $\quad(\alpha|0\rangle+\beta|1\rangle)(|00\rangle+|11\rangle) \quad$ (omitting the $1 / \sqrt{ } 2$ factor)

$$
\begin{aligned}
& =\alpha|000\rangle+\alpha|011\rangle+\beta|100\rangle+\beta|111\rangle \\
& =1 / 2(|00\rangle+|11\rangle)(\alpha|0\rangle+\beta|1\rangle) \\
& +1 / 2(|01\rangle+|10\rangle)(\alpha|1\rangle+\beta|0\rangle) \\
& +1 / 2(|00\rangle-|11\rangle)(\alpha|0\rangle-\beta|1\rangle) \\
& +1 / 2(|01\rangle-|10\rangle)(\alpha|1\rangle-\beta|0\rangle)
\end{aligned}
$$

Protocol: Alice measures her two qubits in the Bell basis and sends the result to Bob (who then "corrects" his state)

## What Alice does specifically

Alice applies


$$
\left\{\begin{array}{rll}
1 / 2|00\rangle(\alpha|0\rangle+\beta|1\rangle) & -D= \\
+1 / 2|01\rangle(\alpha|1\rangle+\beta|0\rangle) & -1 / 2|10\rangle(\alpha|0\rangle-\beta|1\rangle) & - \\
+1 / 2|11\rangle(\alpha|1\rangle-\beta|0\rangle) & =\left\{\begin{array}{ll}
(00, \alpha|0\rangle+\beta|1\rangle) & \text { with prob. } 1 / 4 \\
(01, \alpha|1\rangle+\beta|0\rangle) & \text { with prob. } 1 / 4 \\
(10, \alpha|0\rangle-\beta|1\rangle) & \text { with prob. } 1 / 4 \\
(11, \alpha|1\rangle-\beta|0\rangle) & \text { with prob. } 1 / 4
\end{array} ~\right.
\end{array}\right.
$$

to her two qubits, yielding:

Then Alice sends her two classical bits to Bob, who then adjusts his qubit to be $\alpha|0\rangle+\beta|1\rangle$ whatever case occurs

## Bob's adjustment procedure

Bob receives two classical bits $a, b$ from Alice, and: if $b=1$ he applies $X$ to qubit if $a=1$ he applies $Z$ to qubit

$$
X=\left[\begin{array}{ll}
0 & 1 \\
1 & 0
\end{array}\right] \quad Z=\left[\begin{array}{rr}
1 & 0 \\
0 & -1
\end{array}\right]
$$

yielding: $\begin{cases}00, & \alpha|0\rangle+\beta|1\rangle\end{cases}$

Note that Bob acquires the correct state in each case

## Summary of teleportation



Quantum circuit exercise: try to work through the details of the analysis of this teleportation protocol

## No-cloning theorem

## Classical information can be copied



What about quantum information?


## Candidate:


works fine for $|\psi\rangle=|0\rangle$ and $|\psi\rangle=|1\rangle$
... but it fails for $|\psi\rangle=(1 / \sqrt{ } 2)(|0\rangle+|1\rangle) \ldots$
... where it yields output $(1 / \sqrt{ } 2)(|00\rangle+|11\rangle)$
instead of $|\psi\rangle|\psi\rangle=(1 / 4)(|00\rangle+|01\rangle+|10\rangle+|11\rangle)$

## No-cloning theorem

Theorem: there is no valid quantum operation that maps an arbitrary state $|\psi\rangle$ to $|\psi\rangle|\psi\rangle$

## Proof:



Let $|\psi\rangle$ and $\left|\psi^{\prime}\right\rangle$ be two input states, yielding outputs $|\psi\rangle|\psi\rangle|\mathrm{g}\rangle$ and $\left|\psi^{\prime}\right\rangle\left|\psi^{\prime}\right\rangle\left|g^{\prime}\right\rangle$ respectively

Since $U$ preserves inner products:

$$
\begin{aligned}
& \left\langle\psi \mid \psi^{\prime}\right\rangle=\left\langle\psi \mid \psi^{\prime}\right\rangle\left\langle\psi \mid \psi^{\prime}\right\rangle\left\langle\mathrm{g} \mid \mathrm{g}^{\prime}\right\rangle \text { so } \\
& \left\langle\psi \mid \psi^{\prime}\right\rangle\left(1-\left\langle\psi \mid \psi^{\prime}\right\rangle\left\langle\mathrm{g} \mid \mathrm{g}^{\prime}\right\rangle\right)=0 \text { so } \\
& \left|\left\langle\psi \mid \psi^{\prime}\right\rangle\right|=0 \text { or } 1
\end{aligned}
$$



