

Introduction to Quantum Information Processing

CS 467 / CS 667

Phys 467 / Phys 767

C&O 481 / C&O 681

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Course web site at:

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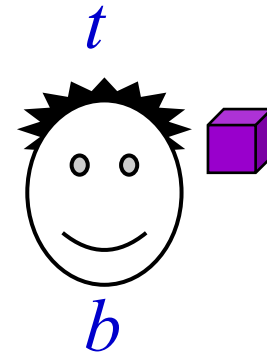
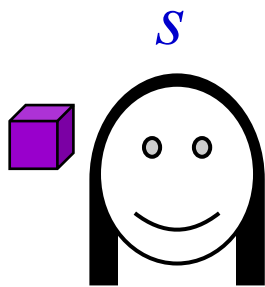
- The Bell inequality and its violation
 - Physicist's perspective
 - Computer Scientist's perspective
- The magic square game
- Communication complexity
 - Equality checking
 - Appointment scheduling (quadratic savings)
 - Are exponential savings possible?
 - Lower bound for the inner product problem
- Simultaneous message passing and fingerprinting

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Bell's Inequality and its violation

Part II: computer scientist's view:

input:



output:

- Rules:**
1. No communication after inputs received
 2. They *win* if $a \oplus b = s \wedge t$



st	$a \oplus b$
00	0
01	0
10	0
11	1

With classical resources, $\Pr[a \oplus b = s \wedge t] \leq 0.75$

But, with prior entanglement state $|00\rangle - |11\rangle$,

$\Pr[a \oplus b = s \wedge t] = \cos^2(\pi/8) = \frac{1}{2} + \frac{1}{4}\sqrt{2} = 0.853\dots$

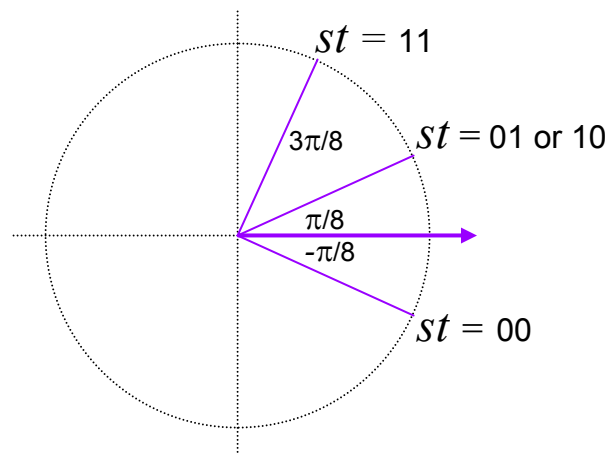
The quantum strategy

- Alice and Bob start with entanglement

$$|\phi\rangle = |00\rangle - |11\rangle$$

- **Alice:** if $s = 0$ then rotate by $\theta_A = -\pi/16$ else rotate by $\theta_A = +3\pi/16$ and measure

- **Bob:** if $t = 0$ then rotate by $\theta_B = -\pi/16$ else rotate by $\theta_B = +3\pi/16$ and measure

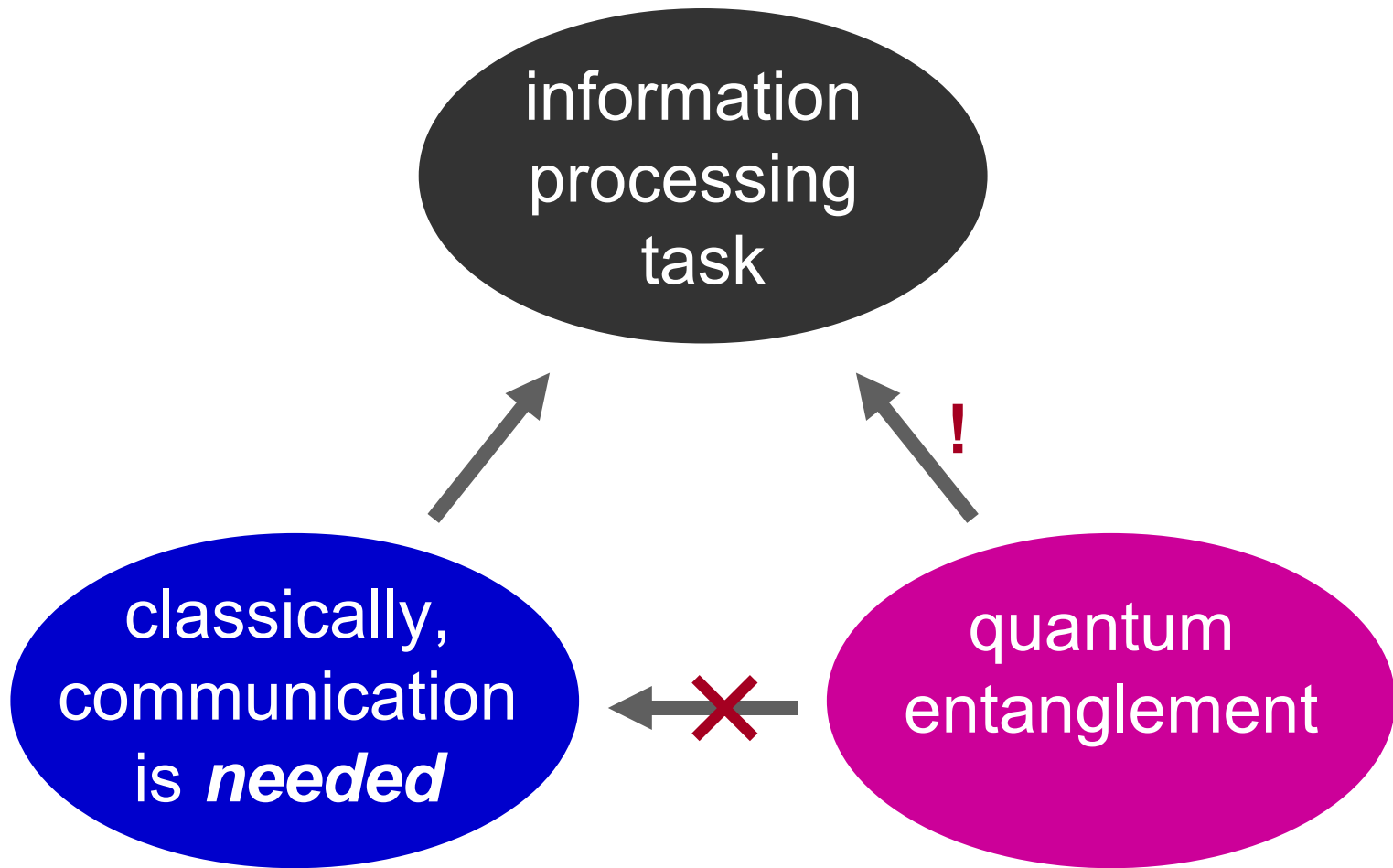


$$\cos(\theta_A - \theta_B) (|00\rangle - |11\rangle) + \sin(\theta_A - \theta_B) (|01\rangle + |10\rangle)$$

Success probability:

$$\Pr[a \oplus b = s \wedge t] = \cos^2(\pi/8) = \frac{1}{2} + \frac{1}{4}\sqrt{2} = 0.853\dots$$

Nonlocality in operational terms



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Magic square game

Problem: fill in the matrix with bits such that each row has even parity and each column has odd parity

a_{11}	a_{12}	a_{13}	even
a_{21}	a_{22}	a_{23}	even
a_{31}	a_{32}	a_{33}	even
odd	odd	odd	

IMPOSSIBLE

		orange
teal	teal	purple
		orange

Game: ask Alice to fill in one row and Bob to fill in one column

They *win* iff parities are correct and bits agree at intersection

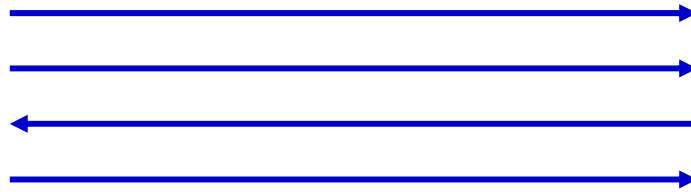
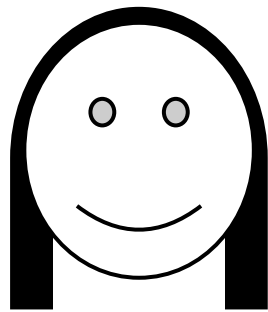
Success probabilities: $8/9$ classical and 1 quantum

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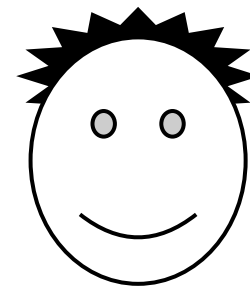
Classical communication complexity

[Yao, 1979]

$x_1 x_2 \dots x_n$



$y_1 y_2 \dots y_n$



$f(x, y)$

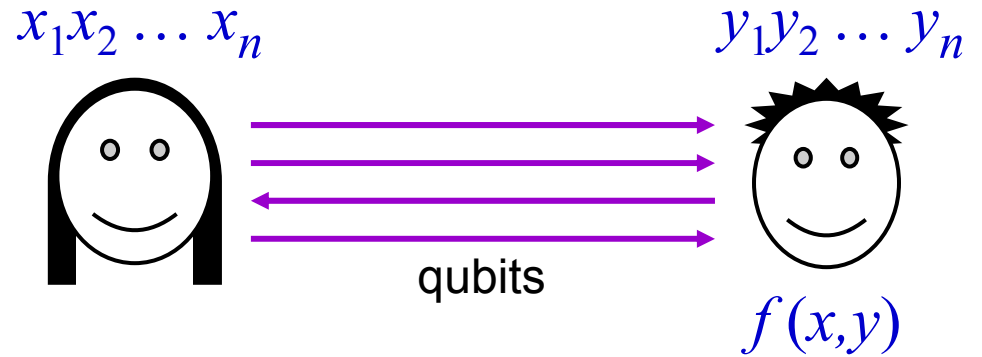
E.g. equality function: $f(x, y) = 1$ if $x = y$, and 0 if $x \neq y$

Any **deterministic** protocol requires n bits communication

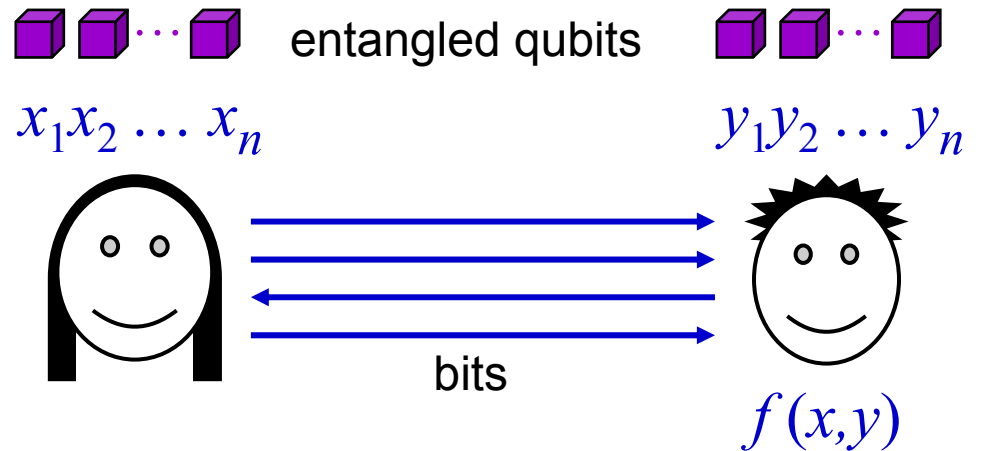
Probabilistic protocols can solve with only $O(\log(n/\varepsilon))$ bits communication (error probability ε), **via random hashing**

Quantum communication complexity

Qubit communication



Prior entanglement

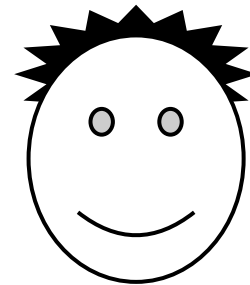
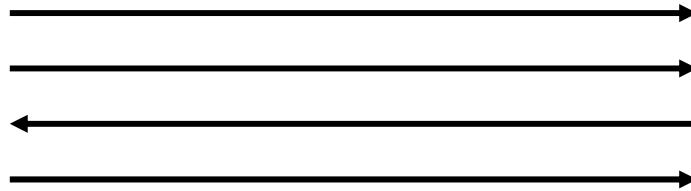
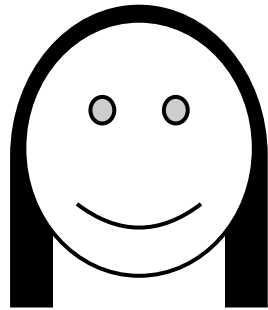


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Appointment scheduling

$$x = \begin{array}{cccccc} 1 & 2 & 3 & 4 & 5 & \dots & n \\ \hline 0 & 1 & 1 & 0 & 1 & \dots & 0 \end{array}$$

$$y = \begin{array}{cccccc} 1 & 2 & 3 & 4 & 5 & \dots & n \\ \hline 1 & 0 & 0 & 1 & 1 & \dots & 1 \end{array}$$



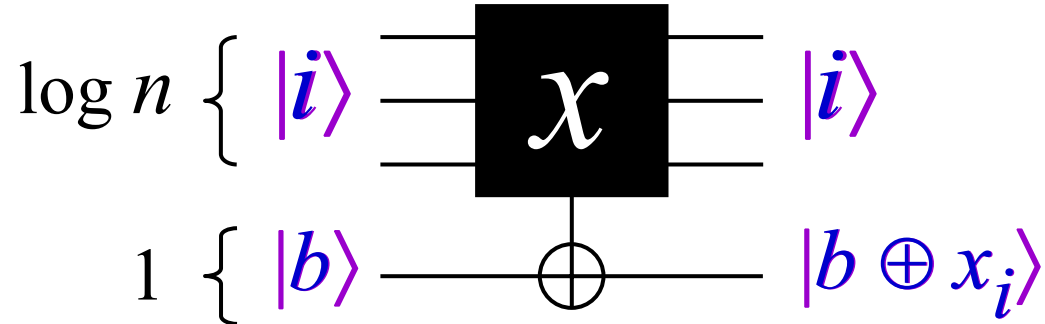
$$i \quad (x_i = y_i = 1)$$

Classically, $\Omega(n)$ **bits** necessary to succeed with prob. $\geq 3/4$

For all $\varepsilon > 0$, $O(n^{1/2} \log n)$ **qubits** sufficient for error prob. $< \varepsilon$

Search problem

Given: $x = \begin{array}{ccccccccc} 1 & 2 & 3 & 4 & 5 & 6 & \dots & n \\ \hline 0 & 0 & 0 & 0 & 1 & 0 & \dots & 1 \end{array}$ accessible via *queries*

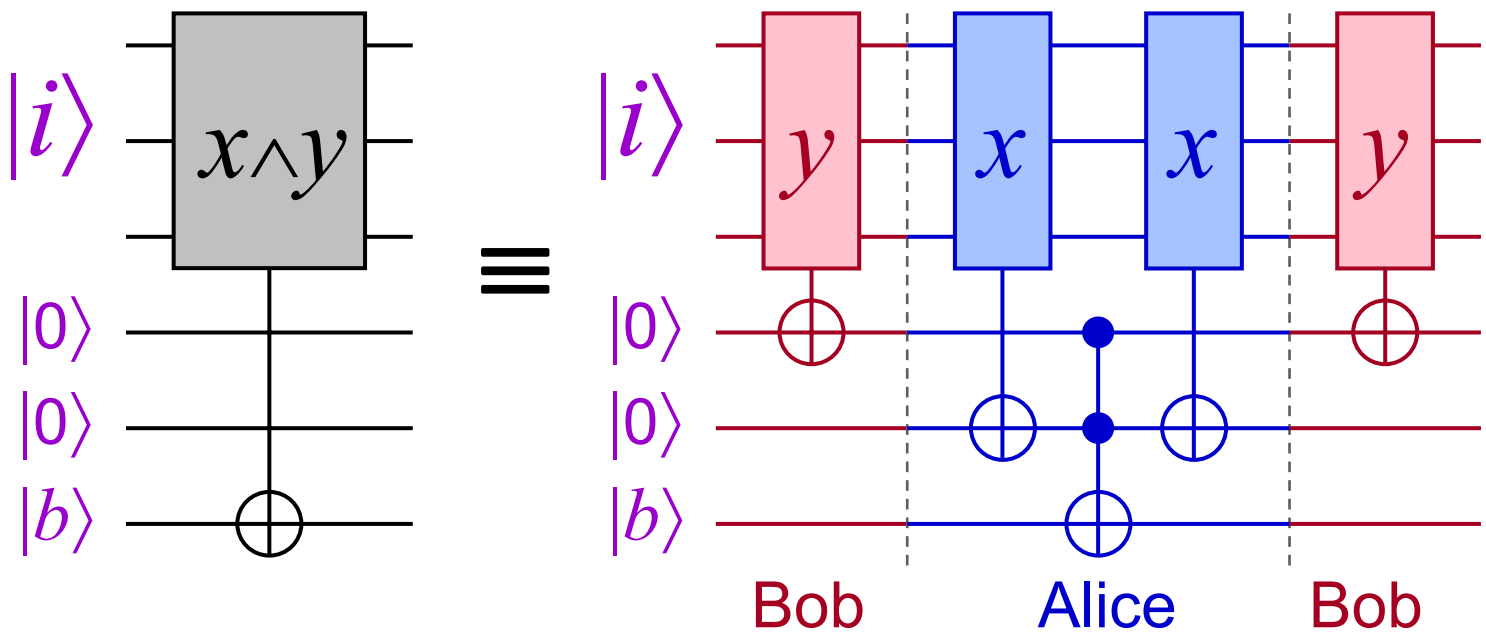


Goal: find $i \in \{1, 2, \dots, n\}$ such that $x_i = 1$

Classically: $\Omega(n)$ queries are necessary

Quantum mechanically: $O(n^{1/2})$ queries are sufficient

		1	2	3	4	5	6	...	n
Alice	$x =$	0	1	1	0	1	0	...	0
Bob	$y =$	1	0	0	1	1	0	...	1
	$x \wedge y =$	0	0	0	0	1	0	...	0



Communication per $x \wedge y$ -query: $2(\log n + 3) = O(\log n)$

Appointment scheduling: epilogue

Bit communication:



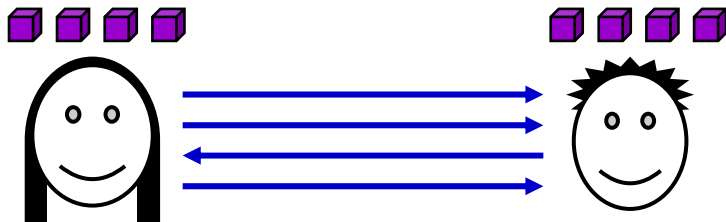
Cost: $\theta(n)$

Qubit communication:



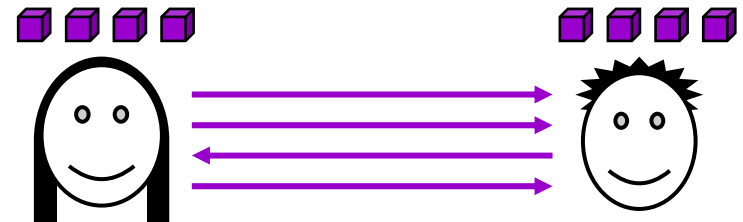
Cost: $\theta(n^{1/2})$ (with refinements)

Bit communication
& prior entanglement:



Cost: $\theta(n^{1/2})$

Qubit communication
& prior entanglement:



Cost: $\theta(n^{1/2})$

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Restricted version of equality

Precondition (i.e. promise): either $x = y$ or $\Delta(x,y) = n/2$

 Hamming distance

(Distributed variant of “constant” vs. “balanced”)

Classically, $\Omega(n)$ bits communication are necessary
for an exact solution

Quantum mechanically, $O(\log n)$ qubits communication
are sufficient ***for an exact solution***

Classical lower bound

Theorem: If $S \subseteq \{0,1\}^n$ has the property that, for all $x, x' \in S$, their *intersection* size is *not* $n/4$ then $|S| < 1.99^n$

Let **some** protocol solve restricted equality with k bits comm.

- 2^k conversations of length k
- approximately $2^n/\sqrt{n}$ input pairs (x, x) , where $\Delta(x) = n/2$

Therefore, $2^n/2^k\sqrt{n}$ input pairs (x, x) that yield **same** conv. C

Define $S = \{x : \Delta(x) = n/2 \text{ and } (x, x) \text{ yields conv. } C\}$

For any $x, x' \in S$, input pair (x, x') **also** yields conversation C

Therefore, $\Delta(x, x') \neq n/2$, implying intersection size is **not** $n/4$

Theorem implies $2^n/2^k\sqrt{n} < 1.99^n$, so $k > 0.007n$

Quantum protocol

For each $x \in \{0,1\}^n$, define $|\Psi_x\rangle = \sum_{j=1}^n (-1)^{x_j} |j\rangle$

Protocol:

1. Alice sends $|\Psi_x\rangle$ to Bob ($\log(n)$ qubits)
2. Bob measures state in a basis that includes $|\Psi_y\rangle$

Correctness of protocol:

If $x = y$ then Bob's result is definitely $|\Psi_y\rangle$

If $\Delta(x,y) = n/2$ then $\langle \Psi_x | \Psi_y \rangle = 0$, so result is definitely **not** $|\Psi_y\rangle$

Question: How much communication if error $1/4$ is permitted?

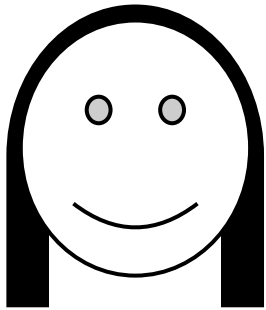
Answer: just 2 bits are sufficient!

Exponential quantum vs. classical separation in bounded-error models

$O(\log n)$ quantum vs. $\Omega(n^{1/4} / \log n)$ classical

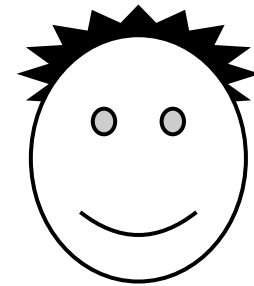
$|\psi\rangle$: a $\log(n)$ -qubit state
(described *classically*)

M : two-outcome measurement



Output: result of
applying M to $U|\psi\rangle$

U : unitary operation
on $\log(n)$ qubits



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Inner product

$$\text{IP}(x, y) = x_1 y_1 + x_2 y_2 + \dots + x_n y_n \pmod{2}$$

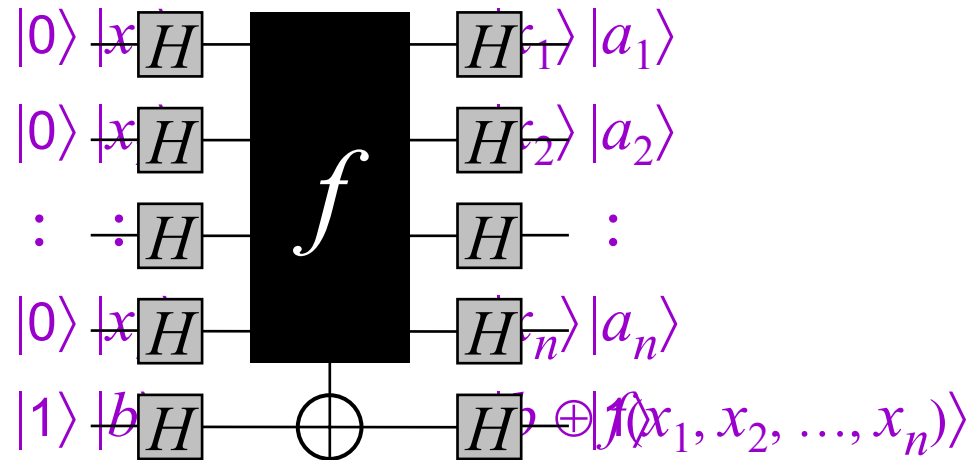
Classically, $\Omega(n)$ bits of communication are required, even for bounded-error protocols

Quantum protocols **also** require $\Omega(n)$ communication

Recall the BV problem

Let $f(x_1, x_2, \dots, x_n) = a_1 x_1 + a_2 x_2 + \dots + a_n x_n \pmod 2$

Given:



Goal: determine a_1, a_2, \dots, a_n

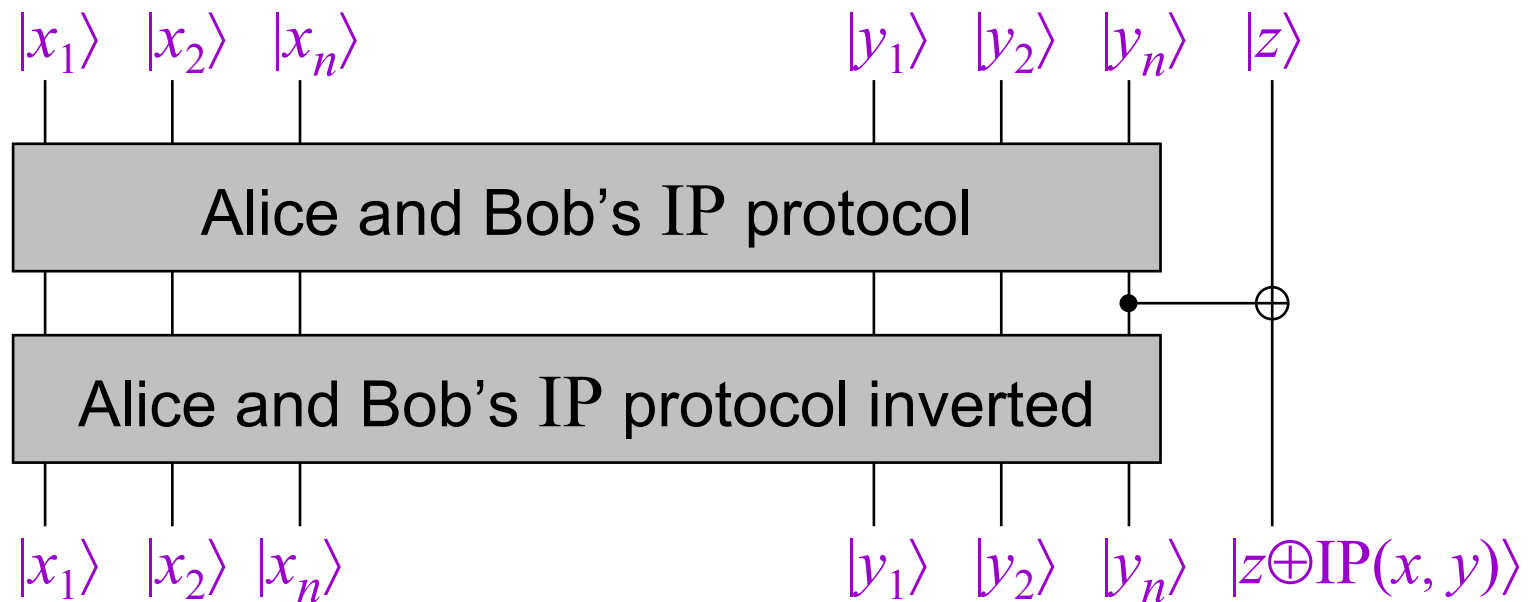
Classically, n queries are necessary

Quantum mechanically, 1 query is sufficient

Lower bound for inner product

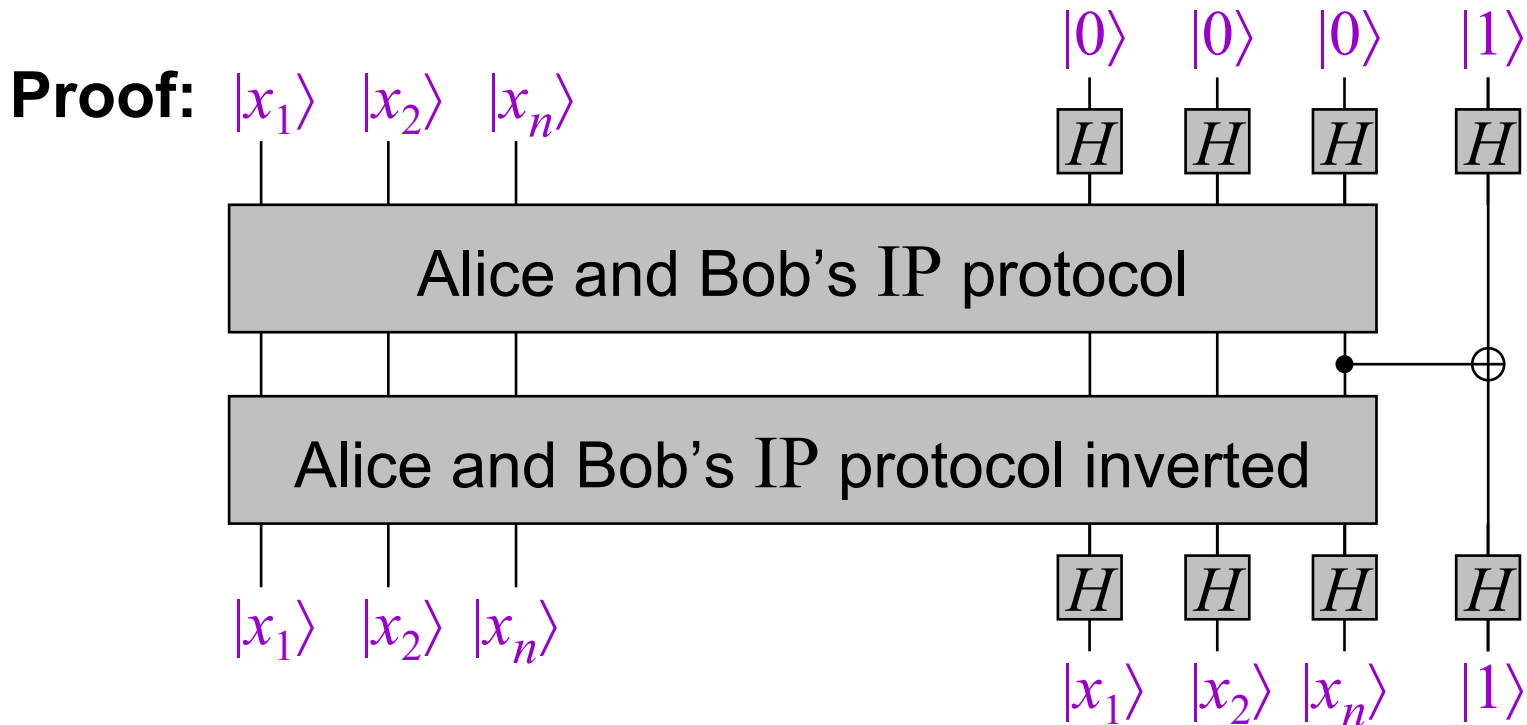
$$\text{IP}(x, y) = x_1 y_1 + x_2 y_2 + \dots + x_n y_n \text{ mod } 2$$

Proof:



Lower bound for inner product

$$\text{IP}(x, y) = x_1 y_1 + x_2 y_2 + \dots + x_n y_n \text{ mod } 2$$



Since n bits are conveyed from Alice to Bob, n qubits communication necessary (by Holevo's Theorem)

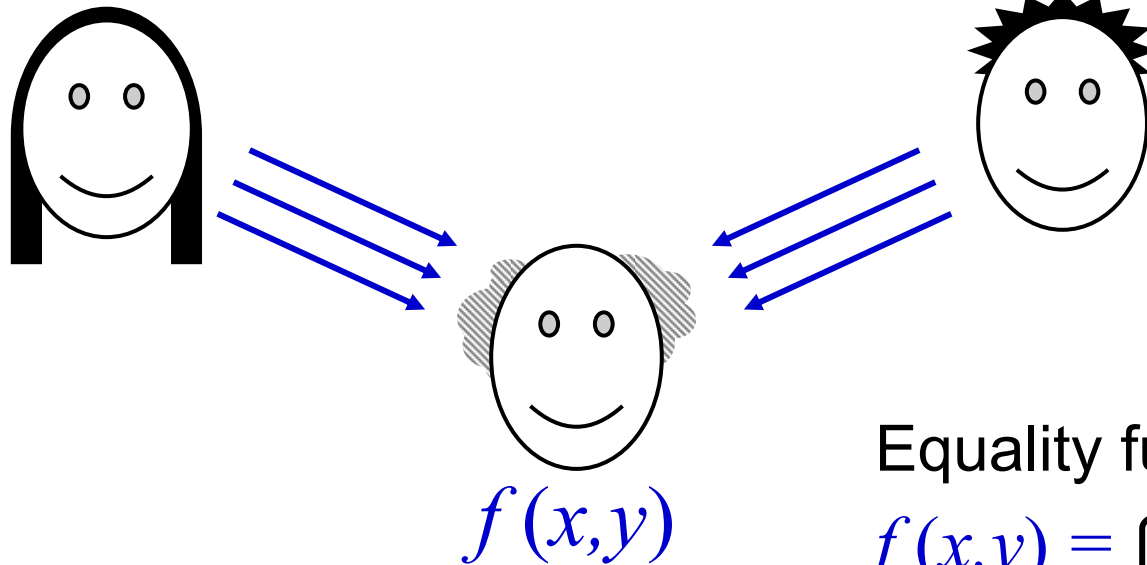
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Equality revisited

in simultaneous message model

$x_1x_2 \dots x_n$

$y_1y_2 \dots y_n$



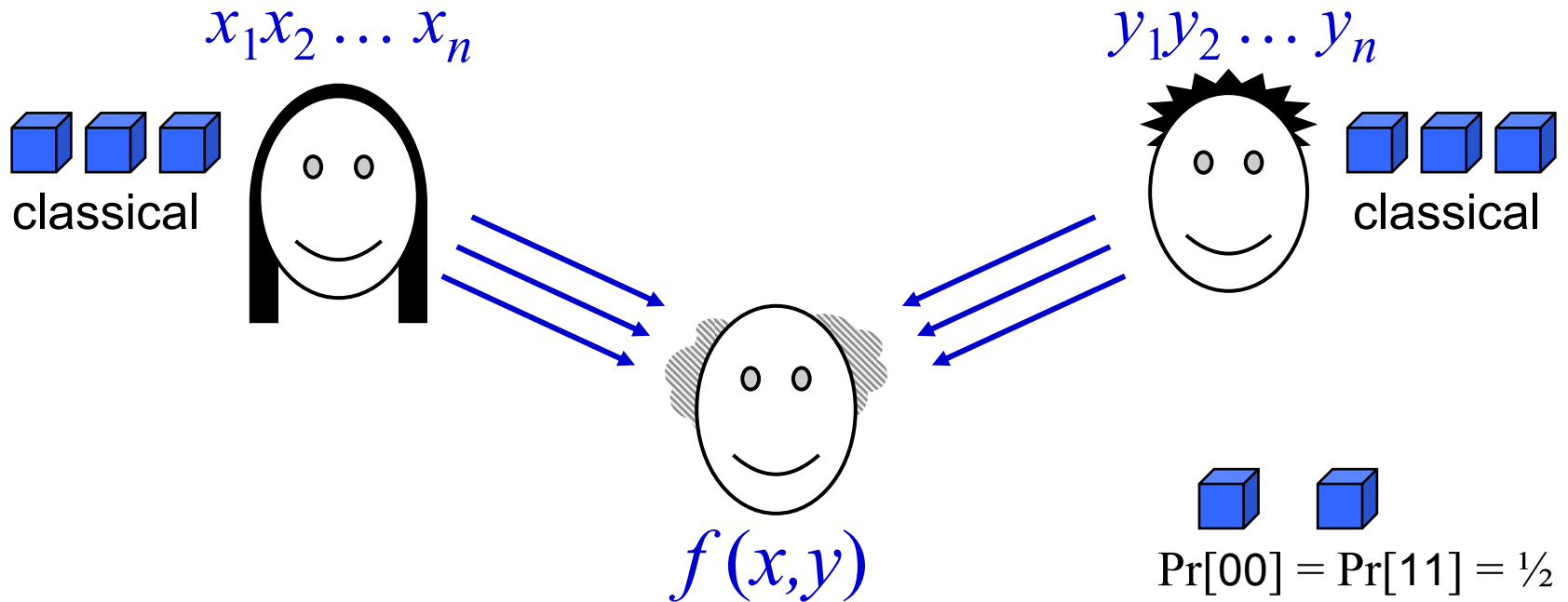
Equality function:

$$f(x, y) = \begin{cases} 1 & \text{if } x = y \\ 0 & \text{if } x \neq y \end{cases}$$

Exact protocols: require $2n$ bits communication

Equality revisited

in simultaneous message model



Bounded-error protocols with a shared random key:
 require only $O(1)$ bits communication

Error-correcting code: $e(x) = 101111010110011001$
 $e(y) = 011010010011001010$

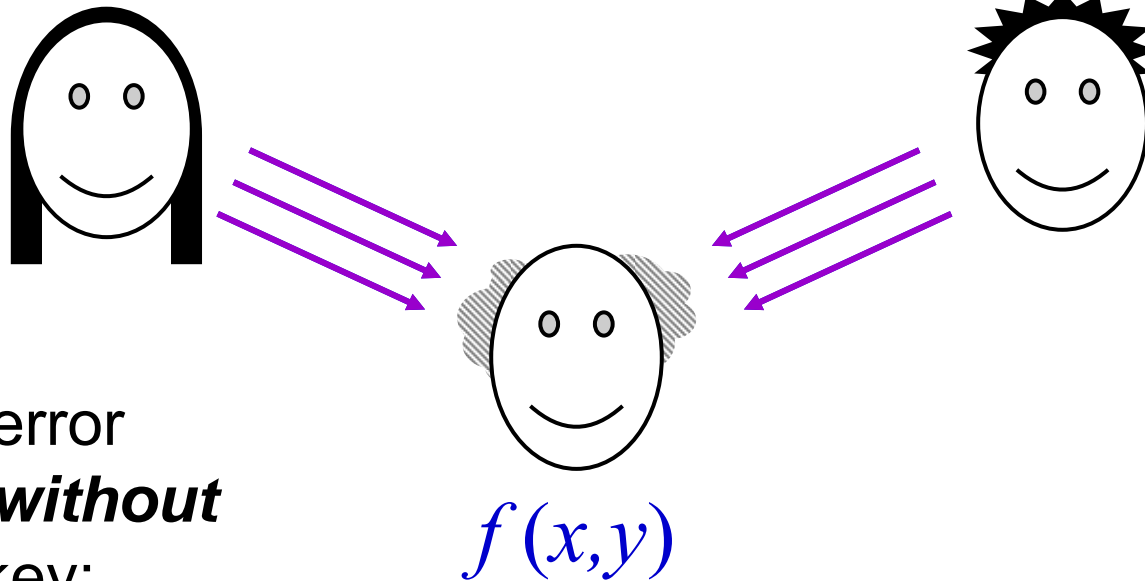
random k

Equality revisited

in simultaneous message model

$x_1 x_2 \dots x_n$

$y_1 y_2 \dots y_n$



Bounded-error
protocols *without*
a shared key:

Classical: $\theta(n^{1/2})$

Quantum: $\theta(\log n)$

Quantum fingerprints

Question 1: how many orthogonal states in m qubits?

Answer: 2^m

Let ε be an arbitrarily small positive constant

Question 2: how many *almost orthogonal** states in m qubits?

(* where $|\langle \Psi_x | \Psi_y \rangle| \leq \varepsilon$)

Answer: 2^{2am} , for some constant $a > 0$

The states can be constructed via a suitable (classical) error-correcting code, which is a function $e: \{0,1\}^n \rightarrow \{0,1\}^{cn}$ where, for all $x \neq y$, $dcn \leq \Delta(e(x), e(y)) \leq (1-d)cn$ (c, d are constants)

Construction of *almost* orthogonal states

Set $|\Psi_x\rangle = \frac{1}{\sqrt{cn}} \sum_{k=1}^{cn} (-1)^{e(x)_k} |k\rangle$ for each $x \in \{0,1\}^n$ ($\log(cn)$ qubits)

Then $\langle \Psi_x | \Psi_y \rangle = \frac{1}{cn} \sum_{k=1}^{cn} (-1)^{[e(x) \oplus e(y)]_k} |k\rangle = 1 - \frac{2\Delta(e(x), e(y))}{cn}$

Since $dcn \leq \Delta(e(x), e(y)) \leq (1-d)cn$, we have $|\langle \Psi_x | \Psi_y \rangle| \leq 1 - 2d$

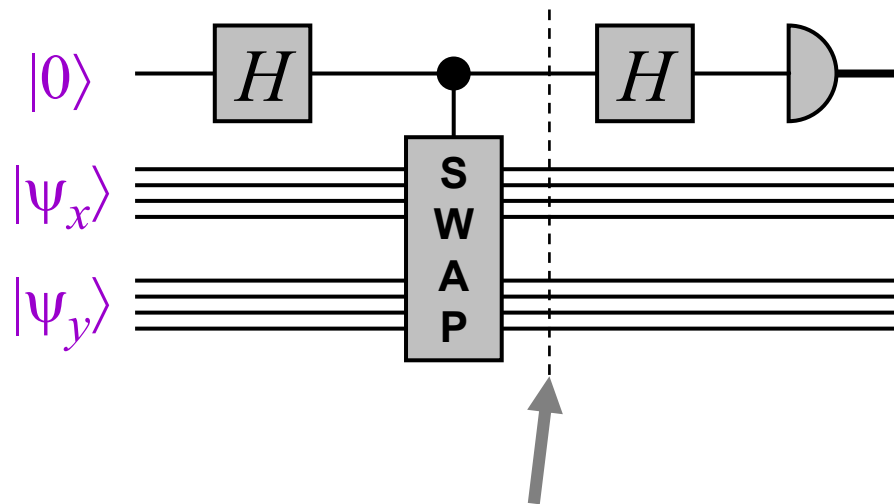
By duplicating each state, $|\Psi_x\rangle \otimes |\Psi_x\rangle \otimes \dots \otimes |\Psi_x\rangle$, the pairwise inner products can be made arbitrarily small: $(1 - 2d)^r \leq \varepsilon$

Result: $m = r \log(cn)$ qubits storing $2^n = 2^{(1/c)2^{m/r}}$ different states

Quantum fingerprints

Let $|\psi_{000}\rangle, |\psi_{001}\rangle, \dots, |\psi_{111}\rangle$ be 2^n states on $O(\log n)$ qubits such that $|\langle \psi_x | \psi_y \rangle| \leq \varepsilon$ for all $x \neq y$

Given $|\psi_x\rangle|\psi_y\rangle$, one can check if $x = y$ or $x \neq y$ as follows:



if $x = y$, $\Pr[\text{output} = 0] = 1$

if $x \neq y$, $\Pr[\text{output} = 0] = (1 + \varepsilon^2)/2$

Intuition: $|0\rangle|\psi_x\rangle|\psi_y\rangle + |1\rangle|\psi_y\rangle|\psi_x\rangle$

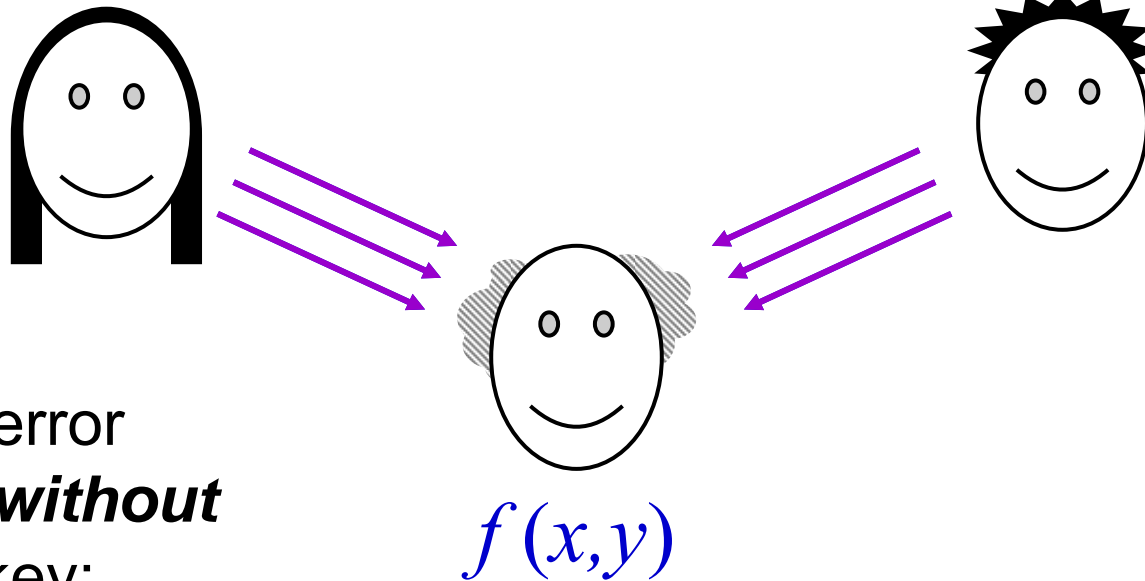
Note: error probability can be reduced to $((1 + \varepsilon^2)/2)^r$

Equality revisited

in simultaneous message model

$x_1 x_2 \dots x_n$

$y_1 y_2 \dots y_n$



Bounded-error
protocols *without*
a shared key:

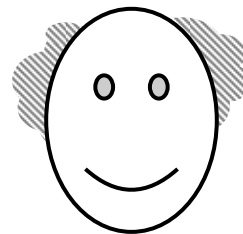
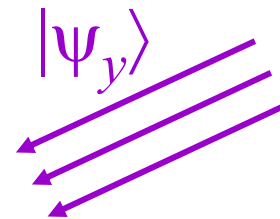
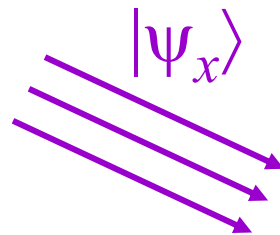
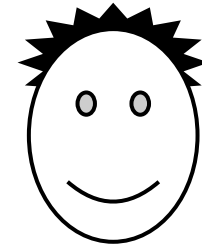
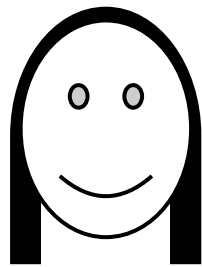
Classical: $\theta(n^{1/2})$

Quantum: $\theta(\log n)$

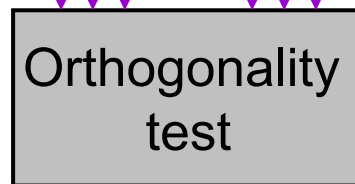
Quantum protocol for equality in simultaneous message model

$x_1 x_2 \dots x_n$

$y_1 y_2 \dots y_n$



$|\Psi_x\rangle$ $|\Psi_y\rangle$



Recall that, *with* a shared key, the problem is easy classically ...

THE END

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